


Simulation of irradiated silicon detectors: an update



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Motivation

- SLHC vs LHC: larger fluences, higher track density, faster readout.
- Silicon detectors operation will require fast, rad-hard, sensitive front-end electronics, smaller pixel size, new (more rad-hard) sensors (DOFZ, Cz, epitaxial, thin silicon, ?)
- Simulation of irradiated silicon detectors allows
 - To connect material properties (N_{eff} and fluence) with detector performance
 - To understand the effect of geometry (electrode size and thickness) on performance
 - To understand if silicon detectors can be operated with high efficiency at fluence around 10^{16} cm^{-1}



Our simulation

From parameterizations of N_{eff} and lifetime as a function of fluence/annealing it computes the performances of irradiated silicon pixel detectors **for different**

- materials
 - StFZ, DOFZ, epi Si, Cz
- doping
 - n+/p, p+/n
- geometries
 - pixel size, thickness
- operating conditions
 - bias, temperature
- front-end electronics parameters
 - threshold, noise



What is new?

In November, we have presented results for StFZ, DOFZ, thin DOFZ detectors with ATLAS pixel size ($50 \times 400 \mu\text{m}^2$)

(Main) upgrades:

- added Cz and epitaxial Si.
- Pixel size is now a job parameter (strip simulation also possible). This talk: smaller pixel size ($70 \times 70 \mu\text{m}^2$) to cope with increased track density at SLHC (use $0.13 \mu\text{m}$ electronics?).
- Pulse time profile and charge collection time.

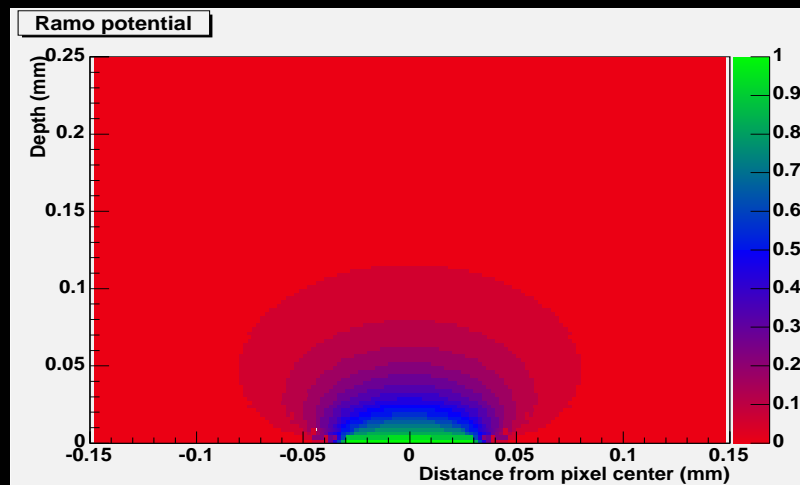
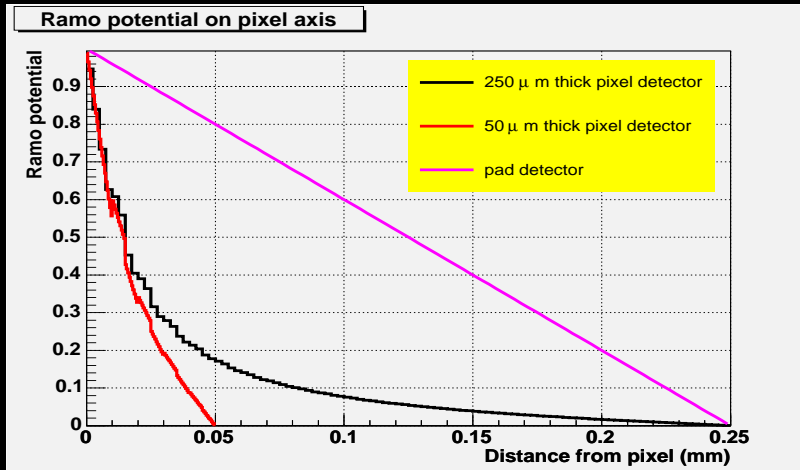


Basics of simulation

- Ionizing particles interactions in the sensors simulated with **Geant4**
- **Charge drift** in silicon (drift, diffusion, trapping).
- Signal induced on pixel electrodes with **Ramo** potential
- **Front-end electronics** response (threshold, noise)
- See the presentation of November for more detailed info.



Small pixels vs pad diode

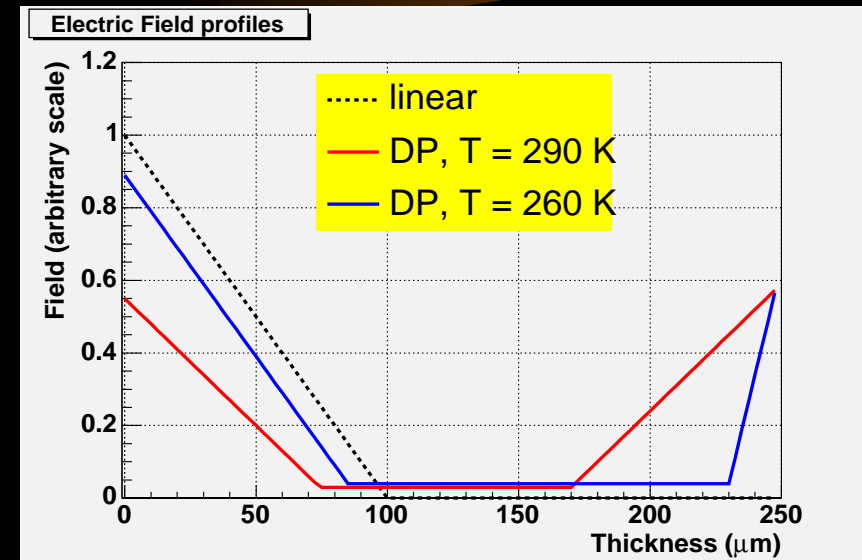


- In a detector with small electrodes most of the signal comes from charges moving near the electrodes.
 - Example: in a 250 μ m thick detector with 50 μ m depletion a charge traversing the depleted region would give **80%** CCE on the nearest pixel. The response of a pad detector (= sum of **negative and positive** signals on all pixels) is only **20%** of the charge!
- In this talk the charge collected is the sum of **positive** signals (because of electronics threshold, negative signals are useless).
Can be **very different** from pad diode CCE.



Electric field

- The field distribution in irradiated silicon has a double peak structure and is a function of dose and temperature. [1-3]
- Power consumption and noise issues require operation at low (about -10°C) temperature to control leakage current.
- At these temperatures the linear field approximation is good for small strips/pixels: charge drift far from the electrodes contributes very little to detector response.



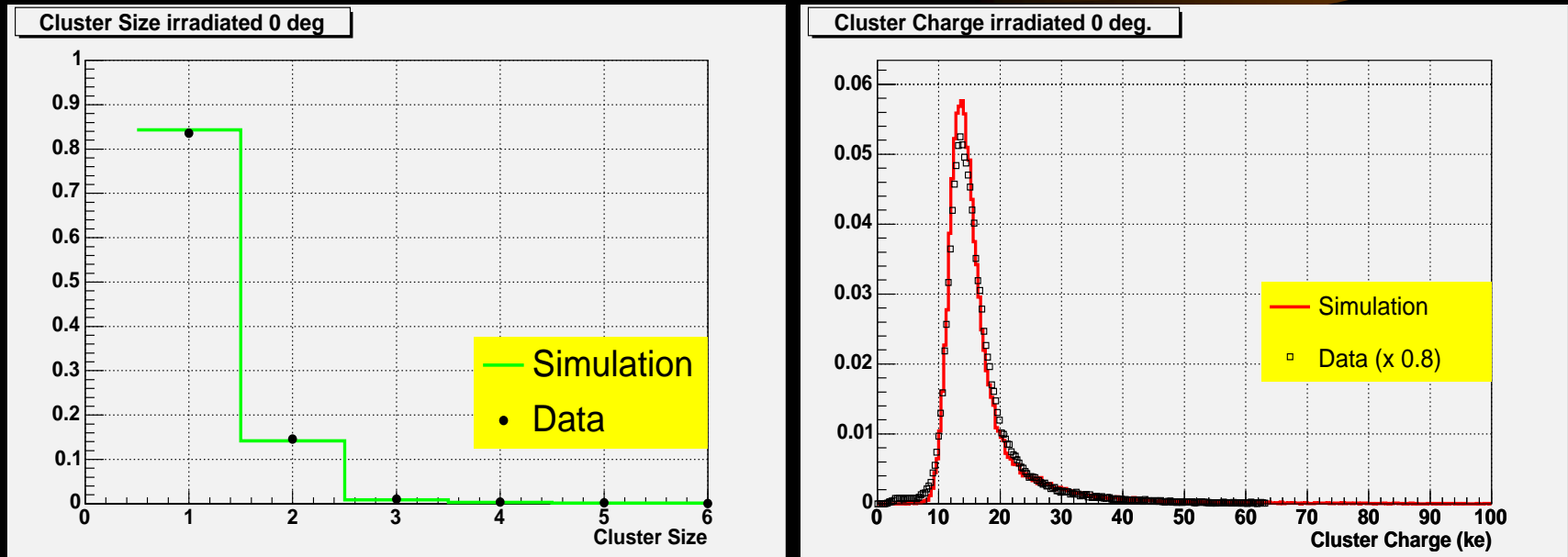
- [1] G. Casse, NIM A426, 140
[2] V. Eremin et al., NIM A360, 458
[3] V. Eremin et al., NIM A476, 556

Presently we use the linear field approximation.



Comparison with data

Simulation validated with experimental data



Comparison with ATLAS Pixel detectors irradiated to $1.1 \cdot 10^{15} n_{eq} \text{ cm}^{-2}$

All parameters from measured values [4-5]

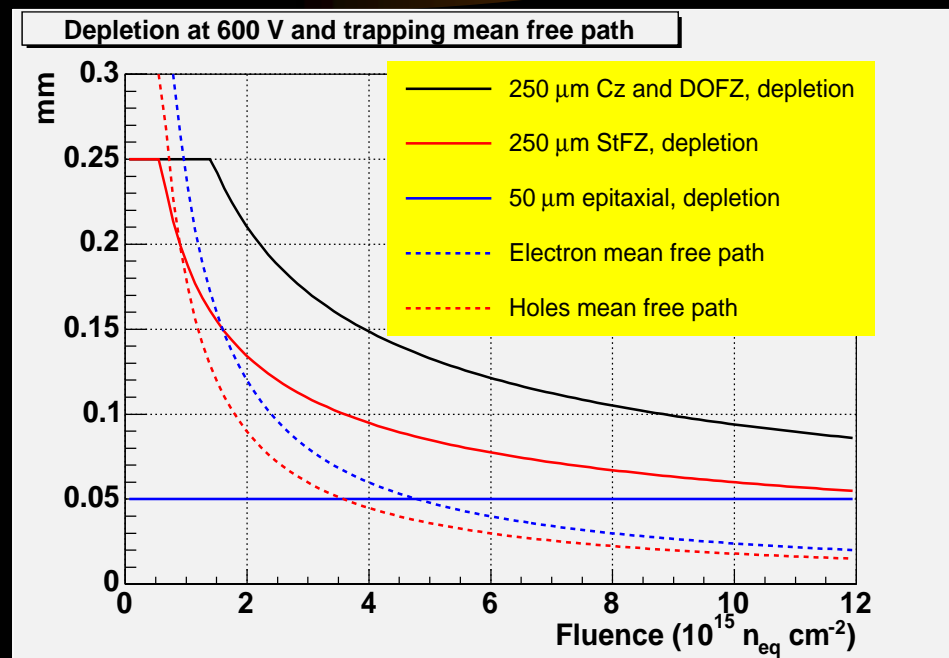
[4] T. Lari, ATL-INDET-2003-015

[5] T. Lari, NIM A518, 349



Radiation damage

- $N_{\text{eff}} = g\Phi$
 $g = 0.023 \text{ cm}^{-1}$ StFZ
 $g = 0.009 \text{ cm}^{-1}$ DOFZ
 $g = -0.009 \text{ cm}^{-1}$ Cz
 $N_{\text{eff}} = -5.79 \cdot 10^{13} \text{ cm}^{-3}$ epitaxial
($V_{\text{fd}} = 100 \text{ V}$ for $50 \mu\text{m}$ sensor)
- $1/\tau = \beta\Phi$
 $\beta_e = \beta_h = 5 \cdot 10^{-16} \text{ cm}^2/\text{ns}$



At high fluences charge collection is limited by trapping mean free path



n-type or p-type?

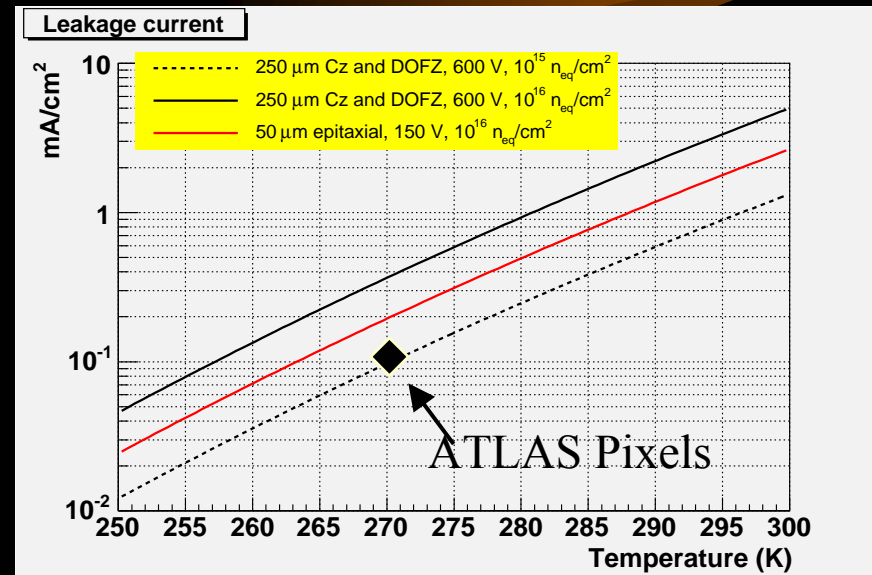
The readout have been chosen to be on the side where the electric field is maximum after irradiation, since this choice results in a better CCE and allows operation in partial depletion mode:

- n-side readout for FZ and DOFZ
- p-side readout for Cz
- p-side readout for epitaxial (first RD50 samples had n-type bulk).



Other parameters

- Thickness 250 μm (as in ATLAS), 50 μm for epitaxial (as first RD50 samples). See November's talk for a DOFZ threshold scan.
- temperature = -10°C
- Bias voltage: 600 V irradiated FZ and Cz (as in ATLAS), 150 V not irradiated and epitaxial.
- Zero incidence angle, no magnetic field



With these parameters leakage current for DOFZ detectors after 10^{16} n_{eq} cm^{-2} is

- 2x smaller (per pixel): less shot noise (35 e for 10 ns integration time)
- 2x larger (per unit area): more power consumption than for ATLAS pixels after 10^{15} n_{eq} cm^{-2} (larger fluence compensated by smaller active volume and temperature)



Charge Collection Vs Fluence

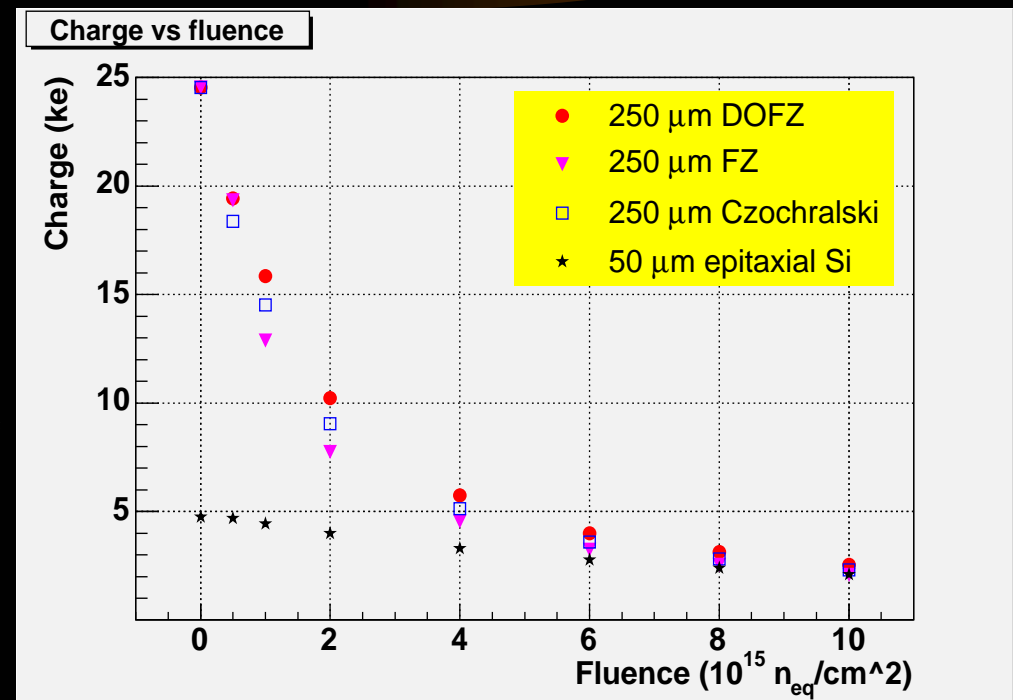
600 V

10^{15} fluence:

- DOFZ better than StFZ when the latter is no longer fully depleted at 600 V
- DOFZ slightly better than Cz (because of n-side signal)
- epitaxial signal very low (because of thin sensor)

10^{16} fluence:

- All detectors are similar (trapping dominant)



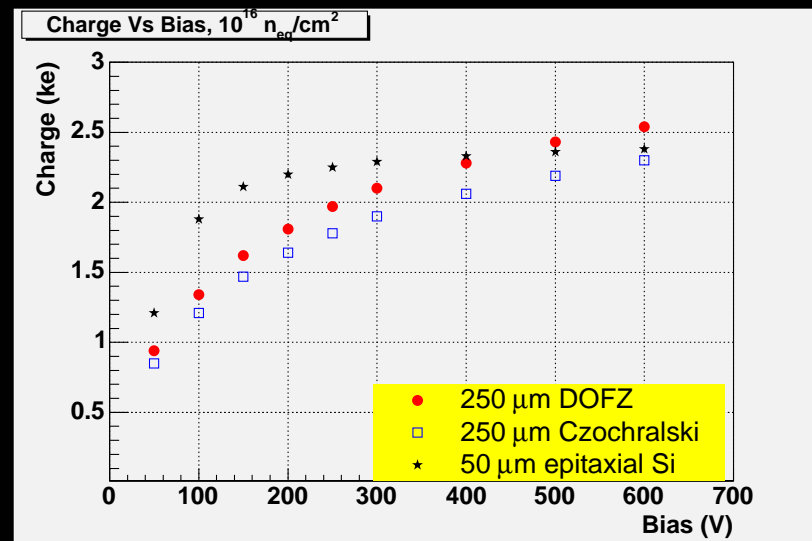
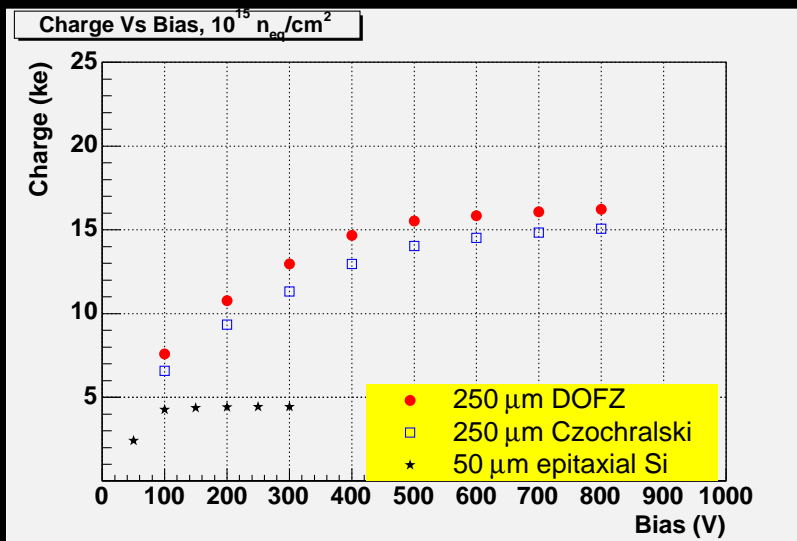
Results should be very similar for strip detectors



Charge collection vs bias voltage

$\Phi = 10^{15} \text{ n/cm}^2$: Cz and DOFZ fully depleted at 440 V, epitaxial at 100 V. Signal increases up to full depletion voltage and is (almost) constant above it.

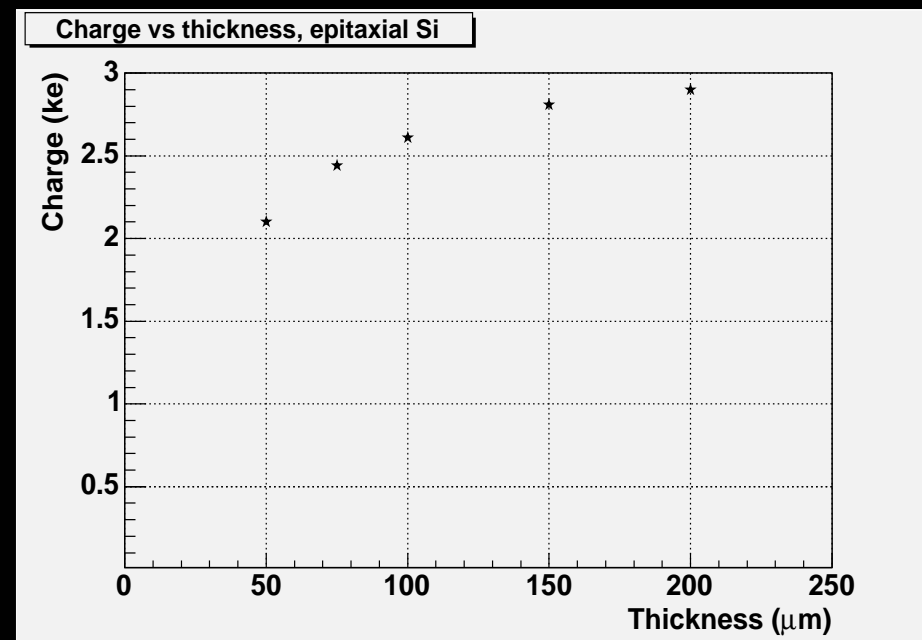
$\Phi = 10^{16} \text{ n/cm}^2$: Cz and DOFZ fully depleted at 4400 V, epitaxial at 100 V. Signal limited by trapping gradually saturates as drift velocity approaches the high-field limit





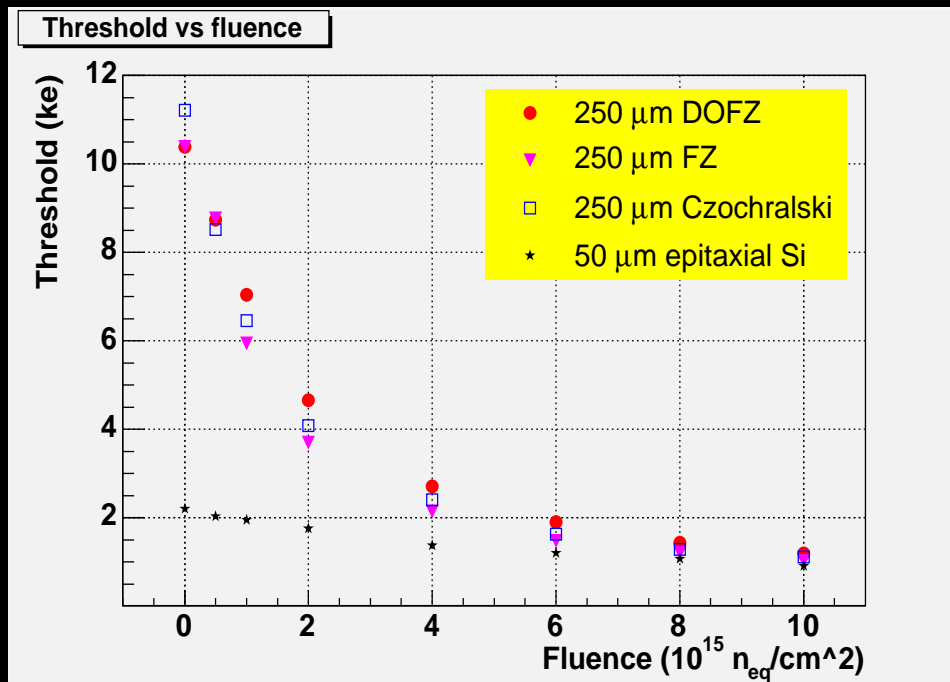
Charge versus epi thickness

- Charge collected after 10^{16} n/cm² as a function of sensor thickness (assuming full depletion)
- First RD50 samples were 50 μ m thick (2100 electrons).
- Asymptotic value is 3000 e⁻, when the thickness is much larger than the mean free path (20 μ m) and the pixel dimensions: charges drifting far from pixels/strip does not contribute to signal.
- Thicker samples are thus predicted to have a significantly larger signal





Threshold and detection efficiency



- The minimum charge which is detected within the trigger window is the **in-time threshold**.
- ATLAS Pixel detectors irradiated to $10^{15} n_{eq} cm^{-2}$ achieve a detection efficiency of 98.2% with an in-time threshold (at 40 MHz) of about 5000 e^- .
- After $10^{16} n_{eq} cm^{-2}$ an in-time threshold of **1000 e^-** is needed (at 80 MHz) to have 97% detection efficiency.

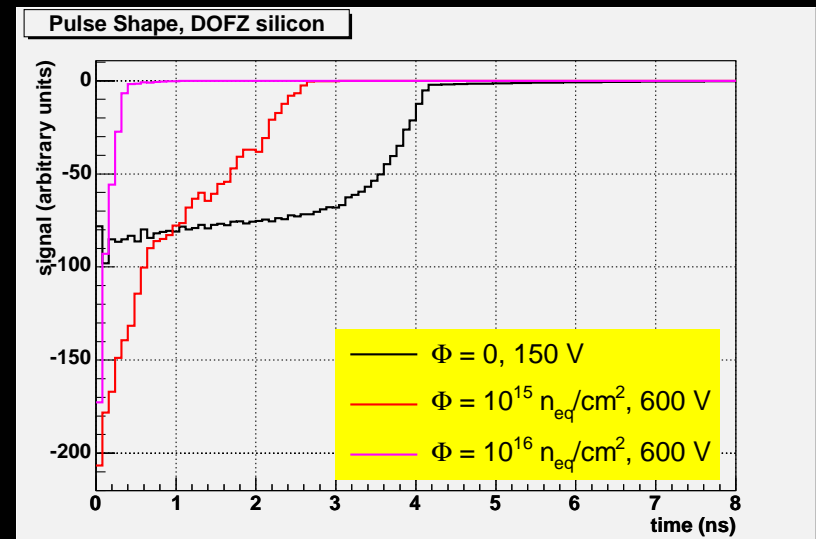
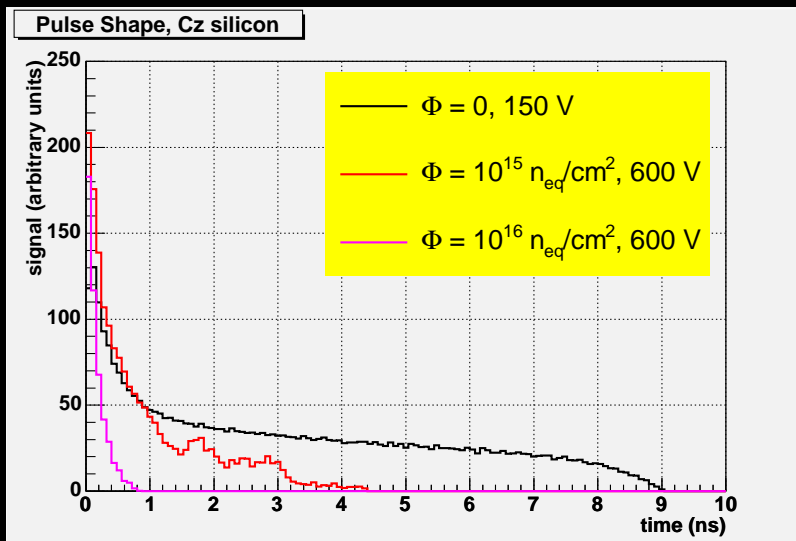
Big challenge for front-end electronics!



Collection time

With stronger irradiation the signal on electrodes is faster (lower depletion, stronger electric field, shorter lifetime).

For thin epitaxial sensors signal is always collected within 1 ns





Conclusions

- The performance of pixel detectors using different silicon materials was simulated after irradiation up to $10^{16} n_{eq} cm^{-2}$
- At the highest fluence, mean signal is **2000-2500 electrons** regardless of material, limited by charge trapping.
- Signal is lower if the thickness is below **100 μm** .
- Sensitivity to **1000 electrons** (fast and low noise rad-hard front-end electronics) is required to operate with high (97%) detection efficiency.